

## THEORY OF IRON DOPED DEUTERATED TRIGLYCINE SULPHATE DOPED TRIGLYCINE SULPHATE (Fe-DTGS) CRYSTAL

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### ABSTRACT

A theoretical approach to explain singular thermodynamic function in order to predict a phase transition in hydrogen bonded ferroelectric Iron doped deuterated triglycine sulphate (Fe-DTGS) crystal and energy propagation in terms of correlation function of inter and intra chain interaction harmonic and anharmonic. This study is extended to third, fourth and higher order phonon anharmonic interactions of two sub-lattice form of the crystal vibrations. Proton tunnelling theory and pseudo vibration model is supported by the blinc and chaudhuri's theory about phase transition in hydrogen bonded ferroelectric crystals. Our dynamic approach is very near to cocharan theory. To define spectral density function, we have implemented time dependent thermal Green's function. Heavy atom displacements are almost perpendicular to spontaneous polarization which is solved by fourth particle cluster theory. Our theory defines the pseudospin character of proton particle in molecular field of a monoclinic system. We have calculated soft mode frequency, lattice frequency shift and width, dielectric constant, loss tangent, quality factor, acoustic attenuation, electric conductivity, relaxation time, smooth function and ratio of figure of merits  $M_2/M_1$ ,  $M_3/M_2$ ,  $M_3/M_1$  characteristics of the Iron doped deuterated triglycine sulphate (Fe-DTGS) crystal. We have considered external electric field effects on ferroelectric properties of the crystal also. Our calculated results have a good agreement with experimental data.

**Keywords:** Figure of merits, pseudo-vibrations, correction functions etc.

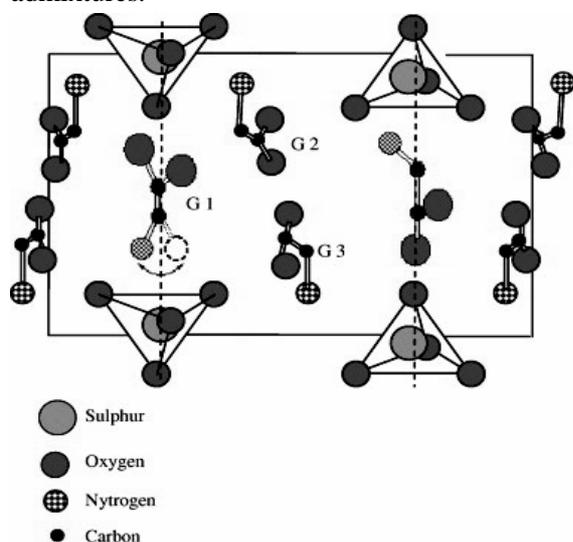
### [I] INTRODUCTION

Triglycine sulphate ( $(\text{NH}_2\text{CH}_2(\text{CH}_3\text{COOH})_3 \cdot \text{H}_2\text{SO}_4)$ ) is a pyroelectric type ferroelectric crystal which is a very famous material in infrared detection and computing technologies etc. has discovered by Matthias et al in 1956 [1]. If its hydrogen atoms (H) is replaced by hydrogen isotopes deuterium (D) then it change in to deuterated triglycine sulphate (DTGS) crystal ( $(\text{ND}_2\text{CD}_2(\text{CD}_3\text{CO OD})_3 \cdot \text{D}_2\text{SO}_4)$ ). Deuterated

triglycine sulphate crystal is monoclinic in both phases. DTGS crystal transform ferorlectric phase  $P_m$  to paraelectric phase  $P_{2/m}$  at 333.86 K. Lattice dimensions of deuterated triglycine sulphate (DTGS) crystal are;  $a=9.462 \text{ \AA}$ ,  $b=12.706 \text{ \AA}$ ,  $c= 5.746 \text{ \AA}$ ,  $V= 690.81 \text{ \AA}^3$  and  $\beta=110.35^\circ$  [2].

Strukov and colleagues have [3–7] examined spontaneous polarization, thermodynamic, calorific relation and electro-conductive

properties of triglycine family crystal (TGS, TGSe). As we see below, that the important consideration is the dependency of entropy  $S$  ( $E$ ,  $T$ ) upon temperature and applied electric field. In this respect it may be useful to study the ferroelectric known to exhibit maximum entropy change at  $T_c$ : ammonium sulphate  $[(NH_4)_2SO_4]$ , for which both ammonium ions and sulphate ions disorder at  $T_c$ , given  $17.6 \text{ J (mol}\cdot\text{K)}^{-1}$ , approximately an order of magnitude larger than in most ferroelectrics [8, 9]. A heat cycle consisting of four steps (or legs), of which two are usually isothermal (or sometimes at constant field) and two are adiabatic [10] published a nice study of this material in 2008, the material's ionic conduction is too high to permit the material to withstand a large field, so commercial devices based on this material are very unlikely. Some chemistry research on the materials science of ammonium sulphate should be carried out to reduce its ionic conductivity, perhaps by admixtures.



**Fig. 1.** Glycine ions (I, II, III) in deuterated triglycine sulphate (DTGS) crystal.

The unit cell of triglycine sulphate (TGS) crystal consists of three glycine groups, namely glycine-I (GI), glycine-II (GII), glycine-III (GIII) and a sulphate ion in an asymmetric unit. The structures of its ferroelectric and paraelectric phases were determined using neutron as well as X-ray

diffraction [11, 12] with the space groups being  $P21$  and  $P21/m$  respectively. The two glycine groups  $GII$  and  $GIII$  are quasi-planar. They are arranged nearly perpendicular to the polar  $b$ -axis and are connected together by a strong hydrogen bond  $GIII-HO_3$ . The hydrogen atom in this H-bond has two equivalent positions of  $HO$  ions, which are closer to  $GIII$  and  $GII$  respectively. The  $GI$  lies close to the  $ac$  plane at  $0.25b$  with its two carbon atoms and two oxygen atoms lying almost in the  $ac$  plane at  $0.25b$  whereas the nitrogen is displaced significantly from the plane. The  $NH_3$  group of  $GI$  has two equivalent positions, one on the left ( $L$ ) and the other on the right ( $R$ ) of the above-mentioned  $ac$  plane. For temperatures higher than transition temperature ( $T_c$ ), the hydrogen in the H-bond  $GIII-H$  ( $GII$  and the  $NH_3$  group of  $GI$  are disordered between their equivalent sites) making the  $ac$  plane (at  $0.25b$  ( $M$ )) a statistical mirror plane [13]. A comparison of the average amplitudes of motion of the heavy atoms near  $T_c$  shows Phase transition in hydrogen bonded triglycine family that the nitrogen atom of

$GI$  has a maximum amplitude in the  $b$  direction, which points to a possible driving role for the nitrogen double-well instability in the mechanism of the ferroelectric transition. Also, the insignificant effect of deuteration on the dielectric properties of TGS (e.g.  $T_c$  is increased by only about 3% in TGS as compared to 80% increase in KDP) suggests that the role of hydrogen tunnelling in the H-bond  $GIII-H$ .  $GII$  may be less dominant as compared to the dynamics of  $NH_3$  group of  $GI$ . The structure of the two ferroelectric domains of TGS are a mirror image of each other [13]. When the polarization of a TGS unit cell reverses, the  $NH_3$  group of  $GI$  moves from one of its equivalent positions, say  $L$  to the other  $R$  and the hydrogen in the H-bond  $GIII-H$  ( $GII$  moves along the bond from one of its equivalent positions  $HO$  to the other  $HO$  ions).

The major disadvantage of using ferroelectric crystals as pyroelectrics is the possibility of a

single domain crystal becoming multi-domain (i.e. depoling of the crystals) during the application. One of the effective ways of preventing depoling in TGS is by doping it with alanine ( $\text{NH}_2\text{CHCH}_3\text{COOH}$ ), which is a molecule similar to glycine ( $\text{NH}_2\text{CH}_2\text{COOH}$ ) with the only difference being the presence of  $\text{CH}_3$  group in alanine in place of one of the two hydrogens of glycine. It is found that the addition of alanine to TGS results in a net internal bias field and the crystals containing alanine retain their self-bias even after prolonged heating above the Curie temperature as well as after the application of large reverse bias electric fields.

Pure TGS crystals possess many disadvantages such as (i) high ferroelectric domain mobility at room temperature, (ii) polarization reversal by electrical mechanical and thermal means, (iii) microbial contamination during the growth and (iv) low Curie temperature. In order to overcome this unfavourable for practical application features variety of dopants such as amino acids, organic and inorganic compounds or deuterated forms have been introduced [14-17]. Due to the competent nonlinear optical efficiency of amino acids and its complexes many researchers have reported with inorganic salts and that has proved as a good candidate for nonlinear optical applications [18-20]. Various metallic ion dopants such as  $\text{Mn}^{2+}$ ,  $\text{Ni}^{2+}$ ,  $\text{Fe}^{3+}$ ,  $\text{Cr}^{3+}$ ,  $\text{Cu}^{2+}$  etc., added to modify the optical properties of TGS crystals [21-24]. Rare earth metal ions like La, Ce and Nd were used to modify the morphology and the coercive field values of TGS crystals. Incorporation of inorganic additives into an organic material can alter the various physical properties in order to achieve better performance in opto-electronic devices. The presence of metallic species in the organic compounds generally raises the hardness of the compounds considerably [25]. Recently alkali metal chloride doped TGS have been reported as high dielectric constant and low hardness values [26]. Infrared

application of DTGS crystal has been proposed by many authors [27].

In the case of deuterated analogs of TGS, it is not possible to completely replace the hydrogen with deuterium. There are seventeen hydrogen in each TGS molecule. Among the seventeen, eleven are active (64.7%) except the six which are in contact with carbon atom. These eleven active H can be substituted with deuterium by recrystallization [28]. DTGS and doped crystals grown during this study have shown to possess superior properties such as materials figures of merit, growth yield and ease of handling over pure DTGS crystals, and are, therefore, very useful and attractive for IR sensing device applications. Pure and doped DTGS crystals, respectively. It can be inferred that presence of dopants in the process of crystal growth results in changes in growth rates along [001] and [010] directions and the growth yield. The growth is anisotropic in nature. It is relatively higher along [010] directions compared to [001] direction; thereby, we observed smaller {010} faces as compared with {001} faces in both pure and doped DTGS crystals. [29]

Pyroelectric detecting materials possess a number of characteristics which are of significance when considering their use in a given application. They can be used over a large spectral bandwidth, the only requirement being that the energy be absorbed. They can be used over a wide range of temperatures. They have low power requirements and can operate for long period on battery power, and last, but not least, they are generally low-cost devices. Some of the devices built with crystals of TGS family are (i) Intruder alarms (ii) Fire alarms (iii) Pollution monitoring and gas analysis (iv) Radio meters (v) Laser detectors (vi) Pyroelectric thermal imaging (vii) Pyroelectric vidicon and (viii) Pyroelectric arrays. DTGS is a successful high sensitive detector type used as infrared spectrometers which is used to take measurements in the mid-infrared (mid IR) range. For better quality results in ferroelectric, dielectric, pyroelectric properties of the crystal,

we have chosen impure Iron doped deuterated triglycine sulphate (Fe-DTGS) crystal. Impurity level of iron in the crystal is considered 0.25% molecular weight. This impurity concentration does not make any structural changes in the crystal body.

### [II] THEORY

To solved Iron doped deuterated triglycine sulphate crystal we have applied two sub lattice pseudospin lattice coupled mode model [30] and Ising spin model [31] of ferroelectrics. In application of these models the major problem was to solve inhomogeneous behaviour of eigen values and eigen functions of the crystal which is solved by Green's function theory [32].

Hamiltonian of Iron doped deuterated triglycine sulphate crystal (Fe-DTGS),

$$\begin{aligned}
 H = & -2\Omega \sum_i (S_{1i}^x + S_{2i}^x) - \sum_{ij} J_{ij} [(S_{1i}^z S_{2j}^z) + (S_{2i}^z S_{1j}^z)] \\
 & - \sum_{ij} K_{ij} (S_{1i}^z S_{2j}^z) - 2\mu E \sum_i (S_{1i}^z + S_{2i}^z) + \frac{1}{4} \sum_k \omega_k (A_k A_k^+ + B_k B_k^+) \\
 & - \sum_{ik} V_{ik} S_{1i}^z A_k - \sum_{ik} V_{ik} S_{2i}^z A_k^+ \dots \dots \dots \\
 & + \sum_{k_1 k_2 k_3} V^{(3)}(k_1, k_2, k_3) A_{k_1} A_{k_2} A_{k_3} \\
 & + \sum_{k_1 k_2 k_3 k_4} V^{(4)}(k_1, k_2, k_3, k_4) A_{k_1} A_{k_2} A_{k_3} A_{k_4} \dots \dots \dots \\
 & + \sum V(\text{Iron} - \text{Interactions}) \\
 \dots(1)
 \end{aligned}$$

Where In equation (1) above  $\Omega$  is proton tunneling frequency,  $S^z$  and  $S^x$  are components of pseudospin variable  $J_{ij}$  is interaction between same lattices and  $K_{ij}$  is interaction between different lattices.  $\mu$  is dipole moment of O-H-O bond,  $E$  is external electric field  $V$  is spin lattice interaction and  $A_k$  and  $B_k$  are position and momentum operators  $\omega_k$  is harmonic phonon frequency  $V^{(3)}$  and  $V^{(4)}$  are third and fourth order atomic force constant

Equation have been solved by the Green's function theory

$$\begin{aligned}
 G_{ij}(t-t') &= \langle \langle S_{1i}^z(t); S_{1j}^z(t') \rangle \rangle \\
 &= -i\theta(t-t') \langle \langle [S_{1i}^z(t); S_{1j}^z(t')] \rangle \rangle, \quad \dots(2)
 \end{aligned}$$

in which  $\theta(t-t')$  step function is zero for  $t < t'$  and unity for  $t > t'$ . The angular bracket  $\langle \dots \rangle$  denotes ensemble average over a grand canonical ensemble.

Spontaneous polarization in Fe-DTGS crystal is  $P_s = 2N_\mu \langle S^z \rangle$  ... (3)

In two dimensional form of pseudospin in z-directional sum is not zero before  $T_C$  (Curie temperature).

$$S_{1i}^z + S_{1j}^z \neq 0;$$

And  $S_{1i}^z \neq S_{1j}^z$ ; For  $T < T_C$ ,

But after curie temperature total spontaneous polarization become zero.

$$S_{1i}^z + S_{1j}^z = 0; \quad \text{For } T > T_C,$$

Specially for Fe-DTGS crystal (ferroelectric) both spin become zero.

$$S_{1i}^z = S_{1j}^z = 0;$$

For Fe-DTGS

(ferroelectric crystal) crystal at phase transition,

while  $S_{1i}^z; S_{1j}^z > 0$ ; but  $P_s \rightarrow 0$  not exactly zero. (for second order phase transition)

$$\langle S^x \rangle = \frac{\Omega}{\tilde{\Omega}} \tanh\left(\frac{\beta \tilde{\Omega}}{2}\right)$$

$$\langle S^y \rangle = 0$$

$$\langle S^z \rangle = \frac{(J_0 \langle S^z \rangle + J_0' \langle S^z \rangle^3)}{2\tilde{\Omega}} \tanh\left(\frac{\beta \tilde{\Omega}}{2}\right) \dots(4)$$

These are obtained as following

Spin shift is

$$\begin{aligned}
 \Delta_s(\omega) &= \frac{a^4}{2\Omega(\omega^2 - \tilde{\Omega}^2)} + \frac{b^2 c^2}{4\Omega\tilde{\Omega}} \\
 &+ \frac{V_{ik}^2 N_k a^2}{2\Omega(\omega^2 - \tilde{\Omega}^2)} + \frac{4\mu^2 E^2 a^2}{2\Omega(\omega^2 - \tilde{\Omega}^2)}
 \end{aligned}$$

and higher order spin shift is

$$\Delta_{s-p}(\omega) = \frac{2V_{ik}^2 \langle S_{1i}^x \rangle \omega_k \delta_{kk'}}{\left[ (\omega^2 - \tilde{\Omega}_k^2)^2 + 4\omega_k^2 \Gamma_k^2(\omega) \right]} \dots(5)$$

Spin width is

$$\begin{aligned} \Gamma(\omega) = & \frac{\pi a^4}{4\Omega\tilde{\Omega}} \left[ \delta(\omega - \tilde{\Omega}) - \delta(\omega + \tilde{\Omega}) \right] \\ & + \frac{b^2 c^2}{4\Omega\tilde{\Omega}} \left[ \delta(\omega - \tilde{\Omega}) - \delta(\omega + \tilde{\Omega}) \right] \\ & + \frac{V_{ik}^2 N_k a^2}{4\Omega\tilde{\Omega}} \left[ \delta(\omega - \tilde{\Omega}) - \delta(\omega + \tilde{\Omega}) \right] \\ & + \frac{2\pi\mu^2 E^2 a^2}{4\Omega\tilde{\Omega}} \left[ \delta(\omega - \tilde{\Omega}) - \delta(\omega + \tilde{\Omega}) \right] \end{aligned}$$

and higher order spin phonon width is

$$\Gamma_{s-p}(\omega) = \frac{4V_{ik}^2 \langle S_{li}^x \rangle \omega_k (\omega^2 - \tilde{\omega}_k^2)}{\left[ (\omega^2 - \tilde{\omega}_k^2)^2 + 4\omega_k^2 \Gamma_k^2(\omega) \right]} \quad \dots(6)$$

### 2.1. Dielectric constant and loss tangent

The response of a dielectric crystal to the external electric field is expressed dielectric susceptibility  $\chi$  given as

$$\chi(\omega) = -\lim_{X \rightarrow 0} 2\pi N \mu^2 G_{ij}(\omega + iX) \quad \dots(7)$$

The  $\chi(\omega)$  is related to dielectric constant as

$$\epsilon = 1 + 4\pi\chi \quad \dots(8)$$

With the help of Eq.(7) and (8) one obtain expression for dielectric constant

$$\epsilon_s(\omega) = (-8\pi N \mu^2) \frac{\langle S_1^x \rangle \Omega}{\left[ (\omega^2 - \hat{\Omega}^2)^2 + 4\Omega^2 \Gamma^2 \right]} \quad \dots(9)$$

$\epsilon(\omega) \gg 1$  in the ferroelectric crystal. The power lost in dielectric when exposed to electromagnetic field is conveniently shown as dielectric tangent loss which is expressed as

$$\tan \delta = \frac{\epsilon''}{\epsilon'} \quad \dots(10)$$

By using Eq. (9) and (10) we obtains expression for loss tangent as

$$\tan \delta = -\frac{2\Omega\Gamma(\omega)}{(\omega^2 - \hat{\Omega}^2)} \quad \dots(11)$$

### 2.2. Quality factor

$$Q\text{-Factor} = 1/\tan\delta \quad \dots(12)$$

Where,  $\tan\delta$ =tangent loss

### 2.3. Acoustic attenuation

The acoustic attenuation is given as

$$\alpha = \frac{\Gamma(\omega)}{\nu} \quad \dots(13)$$

Where,  $\Gamma(\omega)$  is width and  $\nu$  is sound velocity.

### 2.4. Ratio of Figure of Merits

We know that dielectric permittivity can be written as

$$\epsilon = \epsilon' \pm i\epsilon'' \quad \dots(14)$$

where,  $\epsilon'$ =real part of permittivity

$\epsilon''$ =imaginary part of permittivity

and magnitude of this complex number will be

$$\epsilon = \sqrt{(\epsilon')^2 + \epsilon''^2} \quad \dots(15)$$

and dielectric loss

$$\tan \delta = \frac{\epsilon''}{\epsilon'} \quad \dots(16)$$

Solve eq.(15) and eq.(16)

We have

$$\epsilon' = \sqrt{\left( \frac{\epsilon^2}{1 + \tan^2 \delta} \right)}$$

and,

$$\epsilon'' = \sqrt{\frac{\epsilon^2 \tan^2 \delta}{1 + \tan^2 \delta}} \quad \dots(17)$$

in Infra-Red detection we can write the figure of merit can be written as

$$M_1 = \frac{\rho}{\epsilon'}$$

High current responsivity

$$M_2 = \frac{\rho}{\sqrt{\epsilon''}}$$

And vidicons application

$$M_3 = \frac{\rho}{\sqrt{\epsilon'}}$$

Then relative figure of merits will be

$$\frac{M_2}{M_1} = \sqrt{\epsilon''}$$

$$\frac{M_3}{M_2} = \sqrt{\frac{\epsilon''}{\epsilon'}}$$

$$\frac{M_3}{M_1} = \frac{\epsilon''}{\sqrt{\epsilon'}}$$

...(18)

**2.5. Electric conductivity**

$$\sigma = \omega \epsilon_0 \epsilon''$$

...(19)

where,  $\sigma$  =electric conductivity  
 $\omega$  =phonon frequency.

**2.6. Relaxation time (minimum)**

Relaxation time (minimum) is,

$$\tau = \frac{\exp\left(-\frac{\pi}{2} \tan \delta\right)}{\omega}$$

...(20)

**2.7. Differentiability**

$$SmoothFunction. = \frac{2\epsilon''}{\pi\epsilon'}$$

...(21)

**[III] RESULTS**

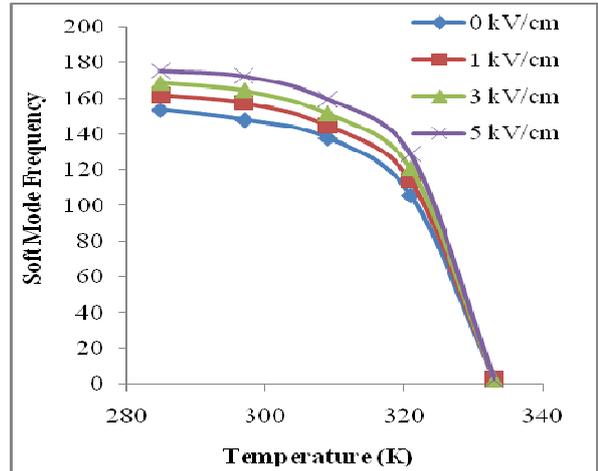
We have calculated soft mode frequency, lattice frequency shift and width, dielectric constant, loss tangent, quality factor, acoustic attenuation, electric conductivity, relaxation time (min.), smooth function and ratio of figure of merits  $M_2/M_1$ ,  $M_3/M_2$ ,  $M_3/M_1$  by using crystal constants given in table 1 and 2. Our all results (fig. 2-14) are in good agreement with experimental data.

**Table: 1.** Crystal constants of deuterated triglycine sulphate crystal.[30]

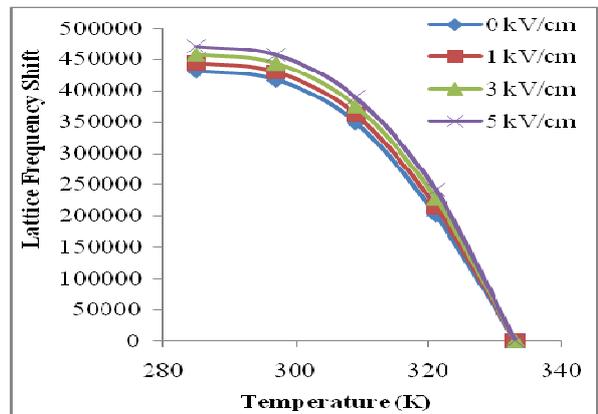
Physical constants	Numerical values for DTGS crystal
$\Omega$ (cm <sup>-1</sup> )	0.01
J (cm <sup>-1</sup> )	470
K (cm <sup>-1</sup> )	0
$V_k$ (cm <sup>-3/2</sup> )	15
$T_c$ (K)	333.86
C (K)	4873.16
$\mu$ (10 <sup>18</sup> esu)	2.3
$\hbar\omega$ (cm <sup>-1</sup> )	1.92

**Table: 2.** Calculated values of  $\langle S_1^x \rangle$ ,  $\langle S_2^x \rangle$ ,  $\langle S_1^z \rangle$ ,  $\langle S_2^z \rangle$  for deuterated triglycine sulphate crystal.

T(K)	$\langle S_1^x \rangle$	$\langle S_2^x \rangle$	$\langle S_1^z \rangle$	$\langle S_2^z \rangle$
285	0.00021	0.0002	0.0025	0.00245
297	0.00043	0.000413	0.0022	0.00221
309	0.00045	0.000445	0.0018	0.00174
321	0.00046	0.00045	0.00125	0.00125
333	0.00047	0.00047	0	0
345	0.00044	0.000434	0	0
357	0.00042	0.000414	0	0
369	0.000425	0.000423	0	0
381	0.000423	0.00042	0	0
393	0.000421	0.000409	0	0



**Fig: 2.** Temperature and electric field dependent soft mode frequency of Fe-DTGS crystal.



**Fig: 3.** Temperature and electric field dependent lattice frequency shift of Fe-DTGS crystal.

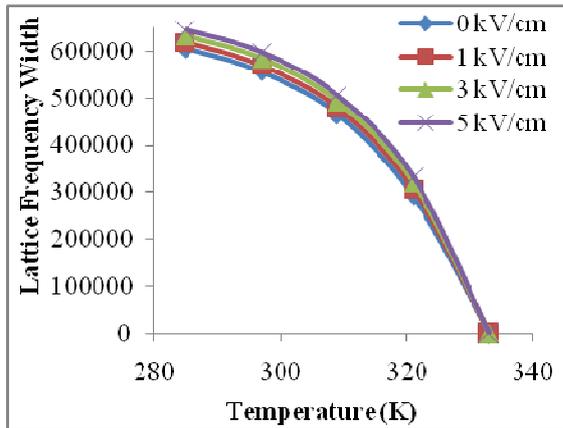


Fig. 4. Temperature and electric field dependent lattice frequency width of Fe-DTGS crystal.

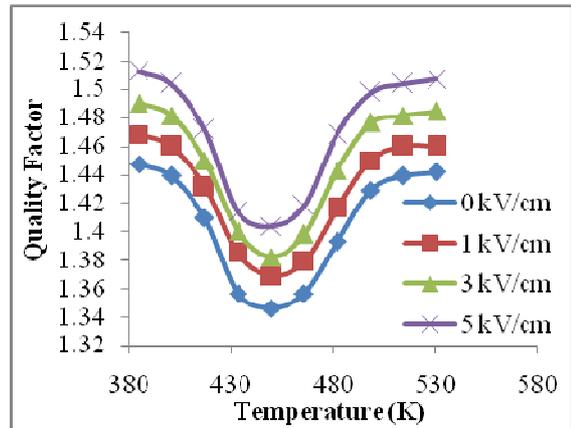


Fig. 7. Temperature and electric field dependent quality factor of Fe-DTGS crystal [33, 34].

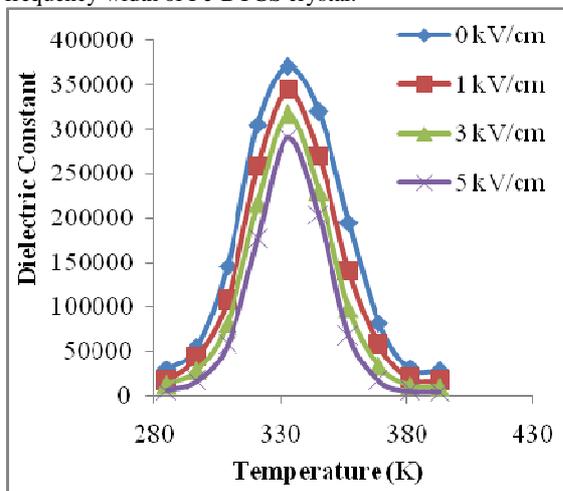


Fig. 5. Temperature and electric field dependent dielectric constant of Fe-DTGS crystal [33, 34].

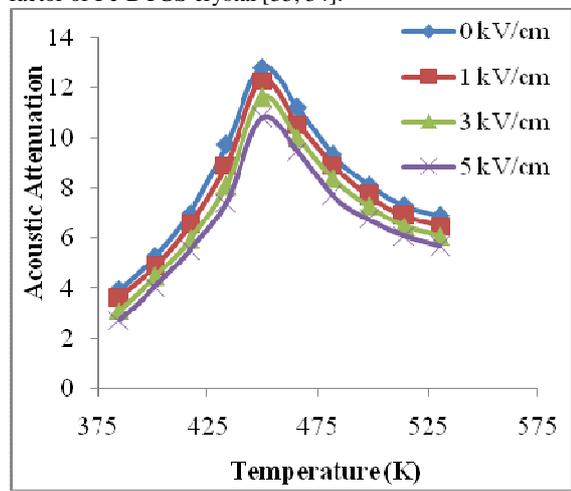


Fig. 8. Temperature and electric field dependent acoustic attenuation of Fe-DTGS crystal [35].

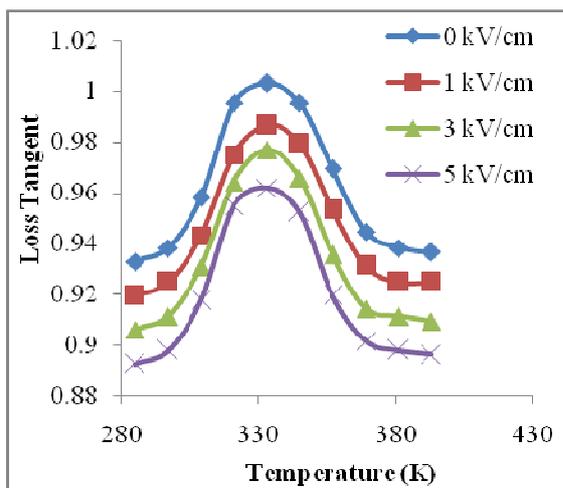


Fig. 6. Temperature and electric field dependent loss tangent of Fe-DTGS crystal [33, 34].

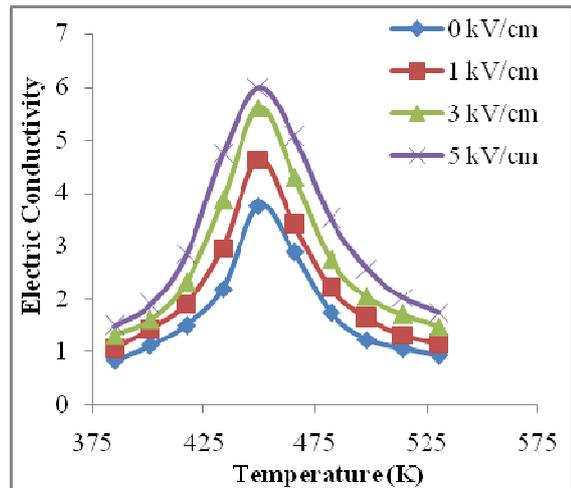
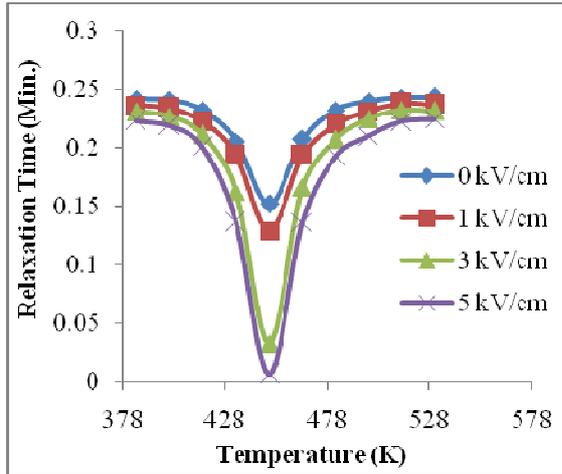
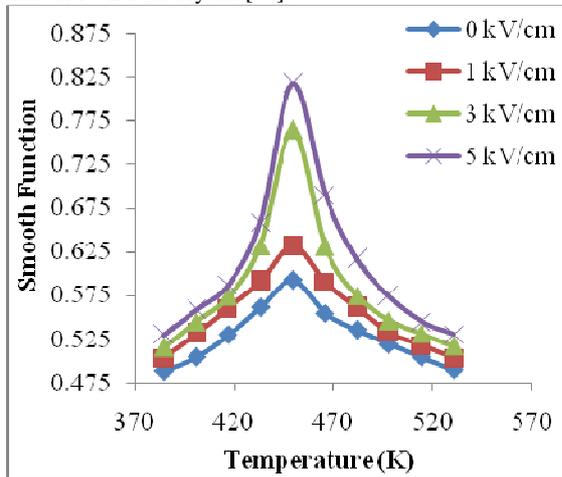


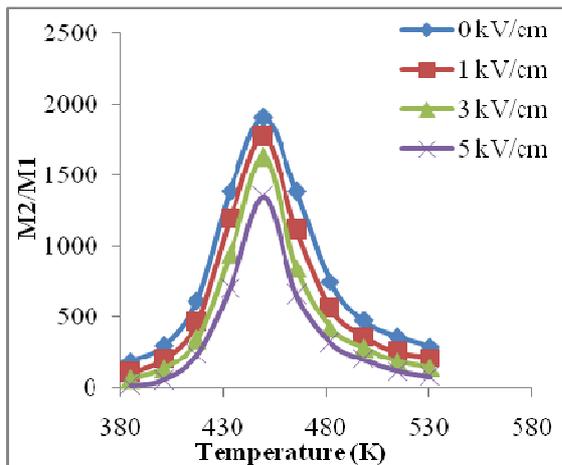
Fig. 9. Temperature and electric field dependent electric conductivity of Fe-DTGS crystal [4, 38].



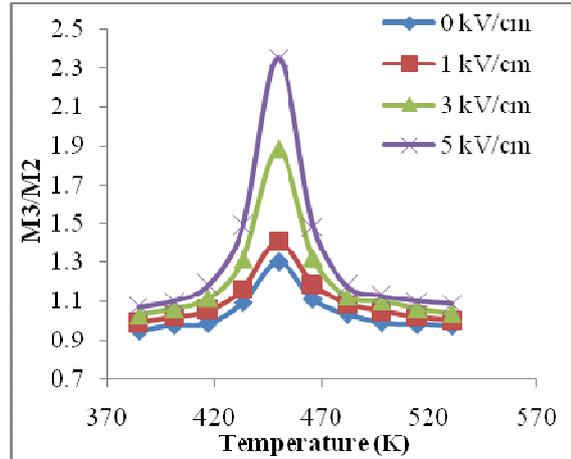
**Fig: 10.** Temperature and electric field dependent relaxation time of Fe-DTGS crystal [36].



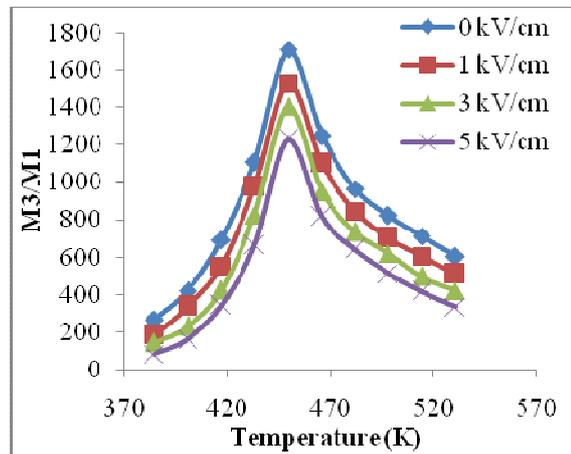
**Fig: 11.** Temperature and electric field dependent smooth function of Fe-DTGS crystal [37].



**Fig: 12.** Temperature and electric field dependent ratio of figure of merits ( $M_2/M_1$ ) of Fe-DTGS crystal [28].



**Fig: 13.** Temperature and electric field dependent ratio of figure of merits ( $M_3/M_2$ ) of Fe-DTGS crystal [28].



**Fig: 14.** Temperature and electric field dependent ratio of figure of merits ( $M_3/M_1$ ) of Fe-DTGS crystal [28].

#### [IV] DISCUSSION

Our calculated results (fig. 2-14) are; soft mode frequency of order  $10^{-2}$ , lattice frequency shift and width of order  $10^{-2}$ , dielectric constant of order  $10^4$  to  $10^5$ , loss tangent of order  $10^0$ , quality factor of order  $10^0$ , acoustic attenuation of order  $10^1$ , electric conductivity of order  $10^0$ , relaxation time (min.) of order of  $10^{-1}$  to  $10^{-2}$ , smooth function of order  $10^{-1}$ , and ratio of figure of merits  $M_2/M_1$  of order  $10^3$ ,  $M_3/M_2$  of order  $10^0$ ,  $M_3/M_1$  of order  $10^3$  at phase transition temperature ( $T_c$ ). When we apply external electric field on the crystal then the corresponding characteristics changes as; soft mode frequency increases, lattice

frequency shift and width increases, dielectric constant decreases, tangent loss decreases, quality factor increases, acoustic attenuation decreases, electric conductivity increases, relaxation time decreases, smooth function increases, and ratio of figure of merits  $M_2/M_1$  decreases,  $M_3/M_2$  increases,  $M_3/M_1$  decreases with increasing order of electric fields. We have observed that increasing electric field slightly decrease the phase transition temperature ( $T_c$ ). Our results show a high current response coefficient, vidicon sensitivity coefficient and infra-red detection coefficient which make this material very sensitive to infrared detection techniques. Fe-DTGS crystal is very useful material in construction of bi-stable ferroelectric memory and quantum state capacitors.

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